# PROPAGATION OF SOLUTIONS OF THE POROUS MEDIUM EQUATION WITH REACTION AND THEIR TRAVELLING WAVE BEHAVIOUR.

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### Index

- Physical motivations. The model
- 2 Key tool. The travelling waves
- Spreading of solutions of compact support
- (4) Spreading of solutions of class  ${\cal A}$
- 5 Convergence of solutions of compact support

### 6 Bibliography

### The Heat Equation.



Porous Medium Equation.



### Physical motivations.

- **Biology** Growth of population depending on its density and a Pearl-Verhaulst type reaction.
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- Astronomy Propagation of intergalatic civilizations.
- Newman, W. I.; Sagan, C. <u>Galactic civilizations: population dynamics</u> and interstellar diffusion. Icarus 46 (1981), 293–327.

## The model.

For m > 1

$$\left\{ \begin{array}{ll} u_t = \Delta u^m + h(u) \quad \text{in } Q := \mathbb{R}^N \times \mathbb{R}_+, \\ u(\cdot, 0) = u_0 \geq 0 \quad \text{in } \mathbb{R}^N, \end{array} \right.$$

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The reaction term h is assumed to be in  $C^1(\mathbb{R}_+)$  and to fulfill, for some  $a \in [0,1)$ ,

$$egin{aligned} &h(0)=0, &h'(1)<0\ &h(u)\leq 0 \mbox{ if } u\in [0,a],\ &h(u)>0 \mbox{ if } u\in (a,1),\ &h(u)<0 \mbox{ if } u>1, \end{aligned}$$

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 h(u) < 0 \text{ if } u > 1, \\
 \int_{0}^{1} m u^{m-1} h(u) \, du > 0
\end{array}$$
(2)

## Arising questions.

- When does the solution *u* spread the value 1 along the medium?
- How fast does it propagate?
- Which shape does the solution take when propagating?

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### The unique travelling wave solution.

If our solution propagates, it probably is similar to a travelling wave solution, a function  $V(x - ct) \equiv V(\xi)$  that satisfies (1), i.e.,

$$\Delta(V^{m}) + cV' + h(V) = 0.$$
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9 / 22

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#### Theorem (Gilding, Kersner).

There exists a minimal speed  $c^* = c^*(m, h) > 0$  such that equation (1) has an unique (up to translations) distinct monotonic "change of fase type" TW solution satisfying

$$\lim_{\xi\to-\infty}V_{c^*}(\xi)=1,\qquad \lim_{\xi\to\infty}V_{c^*}(\xi)=0.$$

and this TW with speed  $c^*$  is  $0 \le V_{c^*} < 1$  and finite.

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For a special pair of functions  $h^*$  and  $\phi$  we consider the following system of ODEs:

$$\begin{cases} f'(t) = h^*(f), \quad f(0) = f_0 > a, \\ g'(t) = c^* \phi(f) - h^*(f)/f, \quad g(0) = g_0, \end{cases}$$

and functions w(x, t) = f(t)V(x - g(t)). Let  $\delta$  be such that h'(u) < 0 for all  $u \in (1 - \delta, 1 + \delta)$ 

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- 1) If  $1 \delta \le f_0 < 1$  then w is a subsolution of (1).
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1) If  $1 - \delta \leq f_0 < 1$  then w is a subsolution of (1).

2) If  $1 + \delta \ge f_0 > 1$  then w is a supersolution of (1).

In both cases,  $w(\xi + c^*t, t) \rightarrow V(\xi - \xi_0)$  and  $(g(t) - c^*t) \rightarrow \xi_0$  when the time t goes to infinity.

# Spreading of solutions of compact support.

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When does the solution *u* spread the value 1 along the medium? How fast does it propagate?

We focus for now on initial data that are non-negative, bounded, piece-wise continuous and compactly supported. We will reffer to this as compactly supported initial data.

### There can always be spreading for every reaction term.

Our first result states that for every reaction h there exist certain initial data of compact support for which our solution propagates. It depends on how much mass  $u_0$  has and how concentrated it is.

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#### Theorem 1.1 (G., Quirós).

There exists a three-parameter  $(x_0, \eta, \rho)$  family of functions v such that if

$$u(x,0) \geq v(x;x,\eta,\rho)$$

for some  $x_0 \in \mathbb{R}^N$ ,  $\eta > 0$  and  $\rho > 0$ , then *u* converges to 1 uniformly on compact sets.

### There can always be spreading for every initial datum.

Next, we see that certain reactions always lead to propagation, regardless on the mass of the initial datum. It depends on the behaviour that h presents near u = 0 compared to the **Fujita exponent**. This is called the **hair-trigger effect**.

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### Theorem 1.2 (G., Quirós).

Suppose that

$$\liminf_{u\to 0}\frac{h(u)}{u^{m+2/N}}>0.$$

and that  $u \neq 0$ . Then *u* converges to 1 uniformly on compact sets.

Fujita exponent for the PME in dimension N:

$$p_F = m + 2/N$$

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# Speed of propagation

If we move from a point  $y_0 \in \mathbb{R}^N$  in a certain direction with a speed c, what will we see as time grows, an empty environment or a saturated one? It depends on said speed.

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#### Theorem 1.3 (G., Quirós).

Whenever spreading happens, for any  $c \in (0, c^*)$ 

$$\lim_{t\to\infty}\min_{|y-y_0|\leq ct}u(y,t)=1.$$

and for any  $c > c^*$ 

$$\lim_{t\to\infty} u(y,t) = 0 \quad \text{ for } \quad |y-y_0| \ge ct.$$

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Too slow will translate to  $c < c^*$  (saturated environment), and too fast to  $c > c^*$  (empty environment). In this sense, the speed  $c^*$  will be called the critical speed of the problem.

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## Spreading of solutions of class $\mathcal{A}$ .

#### Spreading of solutions of class A.

- We focus in dimension N = 1.
- What happens when the initial data are not compactly supported but looks already like a travelling wave?
- Can we get to uniform convergence in the whole space?

## Spreading of solutions of class $\mathcal{A}$ .

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We focus in dimension N = 1.

What happens when the initial data are not compactly supported but looks already like a travelling wave?

Can we get to uniform convergence in the whole space?

We turn our attention to initial data that are non-negative, bounded, piecewise continuous,  $u_0(x) \equiv 0$  for all  $x \ge x_0$ ,  $x_0 \in \mathbb{R}$  and

 $\lim_{x\to-\infty} u_0(x) > \delta,$ 

where  $\delta$  is such that h'(u) < 0 for all  $u \in (1 - \delta, 1 + \delta)$ . This class of functions will be called A class.

## Convergencia uniforme en todo el espacio

We can achieve convergence in all  $\mathbb{R}$  to a travelling wave V solution of (3) with speed  $c^*$ .

#### Theorem 2 (G., Quirós).

Let the dimension be N = 1. Let u be a solution to (1) with  $u_0 \in A$ . Then there exists a  $\xi_0 \in \mathbb{R}$  such that

$$\lim_{t \to \infty} \sup_{x \in \mathbb{R}} |u(x, t) - V(x - c^* t - \xi_0)| = 0,$$

$$\lim_{t \to \infty} \zeta(t) - c^* t = \xi_0.$$
(4)

Here  $\zeta(t)$  denotes the free boundary of the solution.

We present the proof in six steps.

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- 6) We use a stability result, assuring that if our solution is close to a profile V in a certain time, it will remain close for all posterior times, giving convergence along all time sequences.

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Keep in mind that the results about the critical speed  $c^*$  don't say anything about the shape of the function for long times near the front, i.e., near the free boundary.

### Convergence of solutions of compact support in $\mathbb{R}$ .

We can prove, as expected, uniform convergence via a proof similar to the one for the  $\mathcal{A}$  class. We define  $\Omega_1 \equiv [0, \infty)$  and  $\Omega_2 \equiv (-\infty, 0]$ .

#### Theorem 3 (G., Quirós).

Let the dimension be N = 1 and u be a solution of equation (1) of compact support that converges to 1 uniformly in compact sets. Let W be the reflexion of V, i.e.,  $V(\xi) = W(-\xi)$  and  $V_{\xi}$ ,  $W_{\xi}$  translations of said profiles. Then there exist a pair of values  $\xi_*$  and  $\xi^*$  such that

$$\lim_{t \to \infty} |u(x,t) - V_{\xi^*}(x,t)| = 0 \quad \text{in} \quad \Omega_1,$$
 (5)

$$\lim_{t \to \infty} |u(x,t) - W_{\xi_*}(x,t)| = 0 \quad \text{in} \quad \Omega_2.$$
 (6)

## Convergence in greater dimensions.

For radially symmetric initial data of compact support, we expect a Branson correction to appear, i.e.,

$$\lim_{t \to \infty} |u(x,t) - V_{\xi^*}(x,t) + (N-1)c^* \log(t)| = 0.$$

Losely speaking, it comes from the formula for the laplacian in radial coordinates, which transforms our equation in

$$u_t = (u^m)_{rr} + \frac{N-1}{r}(u^m)_r + h(u).$$

# Bibliography

Appart from the already mentioned,

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- Hosono, Y. <u>Traveling wave solutions for some density dependent</u> <u>diffusion equations.</u> Japan J. Appl. Math. 3 (1986), no. 1, 163–196.